# Quantum Computation (CO484) Quantum Physics and Concepts 

Herbert Wiklicky

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7. Gover's Search Algorithm
8. Shor's Quantum Factorisation
9. [Quantum Error Correction]

## Practicalities

Two Lecturers

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Teaching $3 \frac{1}{2}$ weeks until 30 October
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Simon Singh: Code Book, Forth Estate, 1999.

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Manjit Kumar: Quantum - Einstein, Bohr and Their Great Debate about the Nature of Reality, Icon Books 2009

## Photoelectric Effect - Millikan Experiment

## Experimental Setup:

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Experimental Setup:


Observed: The velocity, and thus kinetic energy, of the emitted electrons depends not on the intensity of the incoming light but only on its "colour", i.e. frequency $\nu$.

## Radiation Law

Observed relationship:

$$
W_{k}=h \nu-W_{e}
$$

$W_{k} \ldots$ Kinetic Energy of Electron $W_{e}$... Escape Energy of Material
$\nu$... Frequency of Light
h ... Plank's Constant

$$
\begin{aligned}
& h=6.62559 \cdot 10^{-34} \mathrm{Js} \\
& \hbar=\frac{h}{2 \pi}=1.05449 \cdot 10^{-34} \mathrm{Js}
\end{aligned}
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These were the perhaps most exciting years in the history of theoretical physics, at the same time there were also breakthroughs in special and general relativity, etc.

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In this way one can also explain the spectral emissions (and absorption) of various elements, e.g. to analyse the material composition of stars (and to make great fireworks).

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7. Whereof one cannot speak, thereof one must be silent.

Ludwig Wittgenstein: Tractatus Logico-Philosophicus, 1921

## From Quantum Physics to Computation

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Each area has its own language which however often applies only to classical entities - for the quantum world we often have simply the wrong vocabulary.

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Related Questions: What is our knowledge of what? How do we obtain this information? What is a description on how the system changes?

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Measurement: $m((x, y))=y$, or: $m(\phi)=\sin (\phi)$
Dynamics: $(x, y)(t)=(\cos (t), \sin (t))$ or also: $\phi(t)=t$

## Postulates for Quantum Mechanics [*]

- Observables and states of a system are represented by hermitian (i.e. self-adjoint) elements a of a $\mathrm{C}^{*}$-algebra $\mathcal{A}$ and by states $w$ (i.e. normalised linear functionals) over this algebra.


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- Possible results of measurements of an observable a are given by the spectrum $\mathrm{Sp}(a)$ of an observable. Their probability distribution in a certain state $w$ is given by the probability measure $\mu(w)$ induced by the state $w$ on $\operatorname{Sp}(a)$.


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Key Notions: A quantum systems is (may be) in a certain state, but physicists have to decide which properties they want to observe before a measurement is made (which instrument?).

## Postulates for Quantum Mechanics (ca. 1950) [*]

- The quantum state of a (free) particle is described by a (normalised) complex valued [wave] function:

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\vec{\psi} \in L^{2} \text { i.e. } \int|\vec{\psi}(x)|^{2} d x=1
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- Any observable $A$ is represented by a linear, self-adjoint operator $\mathbf{A}$ on $L^{2}$.
- Possible measurement results are (only) the eigen-values $\lambda_{i}$ of $\mathbf{A}$ corresponding to eigen-vectors/states $\vec{\phi}_{i} \in L^{2}$ with

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- The quantum state of a (free) particle is described by a (normalised) complex valued [wave] function:

$$
\vec{\psi} \in L^{2} \text { i.e. } \int|\vec{\psi}(x)|^{2} d x=1
$$

- Two quantum states can be superimposed, i.e.

$$
\psi=\alpha_{1} \vec{\psi}_{1}+\alpha_{2} \vec{\psi}_{2} \text { with }\left|\alpha_{1}\right|^{2}+\left|\alpha_{2}\right|^{2}=1
$$

- Any observable $A$ is represented by a linear, self-adjoint operator $\mathbf{A}$ on $L^{2}$.
- Possible measurement results are (only) the eigen-values $\lambda_{i}$ of $\mathbf{A}$ corresponding to eigen-vectors/states $\vec{\phi}_{i} \in L^{2}$ with

$$
\mathbf{A} \vec{\phi}_{i}=\lambda_{i} \vec{\phi}_{i}
$$

- Probability to measure (the possible eigenvalue) $\lambda_{n}$ if the system is in the state $\vec{\psi}=\sum_{i} \psi_{i} \vec{\phi}_{i}$ is

$$
\operatorname{Pr}\left(A=\lambda_{n} \mid \vec{\psi}\right)=\left|\psi_{n}\right|^{2}
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There are additional mathematical details in order to deal with "real" quantum physics, e.g. systems an infinite degree of freedom; for quantum computation it is however enough to study finite-dimensional Hilbert spaces.

## Quantum Postulates I - States and Observables

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- An observable is represented mathematically by a selfadjoint matrix (operator) A acting on the Hilbert space $\mathcal{H}$.

Two states can be combined to form a new state $\alpha|\boldsymbol{x}\rangle+\beta|\boldsymbol{y}\rangle$ as long as $|\alpha|^{2}+|\beta|^{2}=1$, by superposition.

Consequence: We can compute with many inputs in parallel.

## Quantum States and Notation

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P.A.M. Dirac "invented" the bra-ket notation (most likely inspired by the limitations of old mechanical type-writers); Simply "take the inner product apart" to denote vectors in $\mathcal{H}$ :

$$
\text { inner product }\langle x \mid y\rangle=\text { product }\langle x| \cdot|y\rangle
$$

For indexed sets of vectors $\left\{\mathbf{x}_{i}\right\}$ (maybe because typographic "typing" was problematic) different notations are used:

$$
\mathbf{x}_{i}=\vec{x}_{i}=\mathfrak{x}_{i}=|i\rangle
$$

## Quantum States and Vectors

Finite quantum states can be described by vectors in $\mathbb{C}^{n}$, e.g.

$$
\vec{\psi}=|\psi\rangle=\binom{1 / \sqrt{2}}{1 / \sqrt{2}}=\frac{1}{\sqrt{2}}\binom{1}{1} \text { or }\langle\phi|=\left(\begin{array}{ll}
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- for enumerating, for example, all eigenvectors of an operator like $\mathbf{A}$ with $|0\rangle=\binom{0}{1}$ and $|1\rangle=\binom{1}{0}$
- to enumerate coordinates of one vector, e.g. $\vec{\psi}_{1}=1 / \sqrt{2}$, or better perhaps: $|0\rangle_{1}=0$.


## Quantum Postulates II - Measurement

- The expected result (average) when measuring observable A of a system in state $|x\rangle \in \mathcal{H}$ is given by:

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\langle\boldsymbol{A}\rangle_{x}=\langle x| \mathbf{A}|x\rangle=\langle x||\mathbf{A} x\rangle
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then we have: $\mathbf{A}=\sum_{i} \lambda_{i} \mathbf{P}_{i}$ (Spectral Theorem).

## Heisenberg's Uncertainty Relation

Theorem
For two observables $\mathbf{A}_{1}$ and $\mathbf{A}_{2}$ we have:

$$
\left(\Delta_{|x\rangle} \mathbf{A}_{1}\right)\left(\Delta_{|x\rangle} \mathbf{A}_{2}\right) \geq \frac{1}{2}\left|\left(\langle x|\left[\mathbf{A}_{1}, \mathbf{A}_{2}\right]|x\rangle\right)\right|
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A standard example of so-called incomensurable observables are position $\mathbf{A}_{1}=x$ and momentum $\mathbf{A}_{2}=p$ (on an infinitedimensional Hilbert Space $\mathcal{H}$ ) for which $[x, p]=i \hbar$ and thus:

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In classical physics observables always commute, are comensurable, i.e. $\left[\mathbf{A}_{1}, \mathbf{A}_{2}\right]=0$. In quantum physics for most observables $\left[\mathbf{A}_{1}, \mathbf{A}_{2}\right] \neq 0$, i.e. the observable algebra is typically non-commutative or non-abelian (cf. multiplication of (complex) numbers vs multiplication of matrices).

## Quantum Dynamics

- The dynamics of a (closed) system is described by the Schrödinger Equation:

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i \hbar \frac{d|x\rangle}{d t}=\mathbf{H}|x\rangle
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Theorem
For any self-adjoint operator $\mathbf{A}$ the operator

$$
\exp (i \mathbf{A})=e^{i \mathbf{A}}=\sum_{n=0}^{\infty} \frac{(i \mathbf{A})^{n}}{n!}
$$

is a unitary operator.

## Irreversible vs Reversible

There are a number of immediate consequence of the postulates.

1. The state develops reversibly, i.e. $\left|x_{t}\right\rangle=\mathbf{U}_{t}\left|x_{0}\right\rangle$ for some unitary matrix (operator).
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The mathematical structure has also consequences for any Quantum Logic, e.g. De Morgan fails, 'Tertium non datur' is not guaranteed, etc.

## Quantum Physics vs Quantum Computation

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Given a quantum system (device). What is its dynamics?

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## Quantum Computation

Given a desired computation (dynamics).
What quantum device (e.g. circuit) is needed to obtain this?

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When will (cheap) quantum computers be available? What will be a killer application for quantum computation? When will we reach quantum supremacy?

